

**Correction to "On the Question of Computation  
of the Dyadic Green's Function at the  
Source Region in Waveguides and Cavities"**

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In the above paper,<sup>1</sup> the following typographical errors should be corrected.

On page 763, the expression following (13) needs a change of

sign and should read as

$$\Gamma_{nm}^2 = -k^2 + \left(\frac{n\pi}{a}\right)^2 + \left(\frac{m\pi}{b}\right)^2.$$

On the same page, the last term in (17) should be

$$\frac{n\pi}{a} \cos \frac{n\pi x}{a} \sin \frac{n\pi x'}{a} \cos \frac{m\pi y}{b} \cos \frac{m\pi y'}{b} \hat{z} \hat{y}.$$

On page 764, the right-hand side of both (26) and (28) should be multiplied by a minus sign.

**ACKNOWLEDGMENT**

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<sup>1</sup> Y. Rahmat-Samii, *IEEE Trans. Microwave Theory Tech.*, vol. MTT-23, pp. 762-765, Sept. 1975.

# Computer Program Descriptions

## Transmission From a Rectangular Waveguide into Half Space Through a Rectangular Aperture

**PURPOSE:** The program calculates the admittance seen by the waveguide, the tangential electric field in the aperture, and the radiation gain patterns for a rectangular waveguide feeding a rectangular aperture in a perfectly conducting plane of infinite extent.

**LANGUAGE:** Fortran IV.

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**AVAILABILITY:** ASIS-NAPS Document No. 03178. Also available from the authors as a report [1]. Card decks are available from the authors at a cost of \$15.00.

**DESCRIPTION:** The problem being considered is illustrated in Fig. 1. It consists of a rectangular waveguide of arbitrary dimensions exciting half space through a rectangular aperture of arbitrary dimensions within the waveguide cross section. The half space is bounded by a perfectly conducting plane of infinite extent.

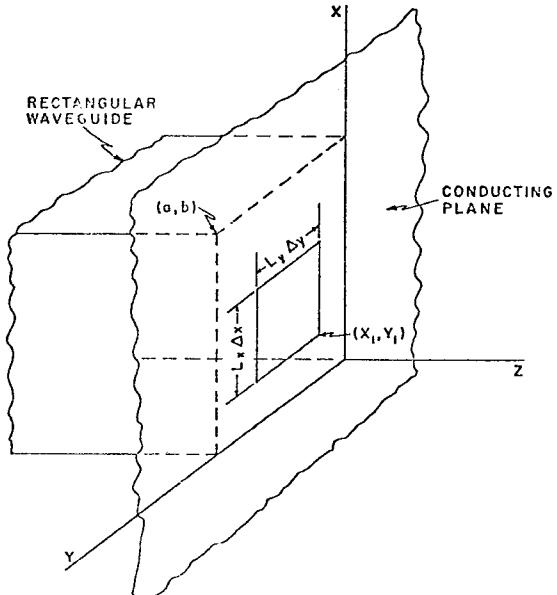


Fig. 1. A rectangular waveguide radiating through a rectangular aperture into half space bounded by an electric conductor.

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The method of solution is that of the generalized network formulation [2]. This consists of dividing the problem into two regions by covering the aperture with an electric conductor and placing the magnetic current  $\mathbf{M} = \mathbf{n} \times \mathbf{E}$  on the waveguide side of this conductor and  $-\mathbf{M}$  on the half-space side of this conductor. By having  $+\mathbf{M}$  on one side and  $-\mathbf{M}$  on the other side

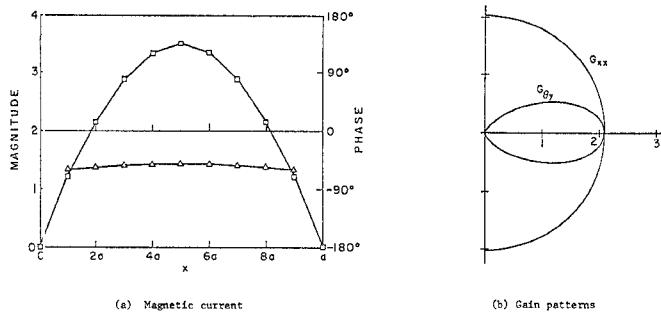


Fig. 2 Equivalent magnetic current  $M_x$  and radiation gain patterns for a rectangular waveguide of dimensions  $\lambda$  by  $\lambda/2$  radiating through a centered rectangular slot of dimensions  $\lambda$  by  $\lambda/10$ . Squares denote magnitude and triangles denote phase.

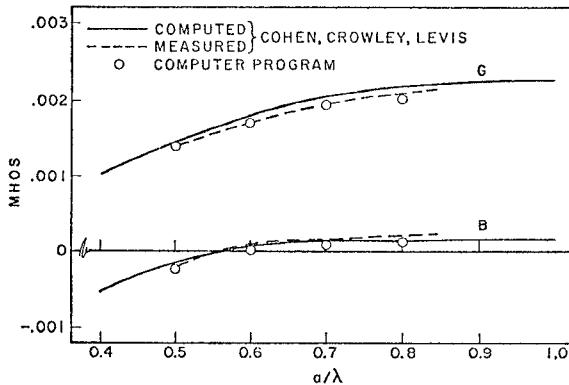


Fig. 3. The equivalent aperture admittance seen by the dominant mode for an open-ended square waveguide of width  $a$  radiating into half space. Our computed results are compared to those calculated and measured by Cohen, Crowley, and Lewis [5].

the boundary condition that tangential  $E$  be continuous across the aperture is always satisfied. Enforcement of the remaining boundary condition that tangential  $H$  be continuous across the aperture leads to an integral equation for the unknown  $M$ . The method of moments [3] is applied to this integral equation to obtain generalized admittance matrices for each region. These generalized admittance matrices are then combined in parallel with current sources dependent on tangential  $H$  from incident field to determine the solution. The expansion and testing functions used in the solution are triangles in the direction of  $M$  and pulses in the direction transverse to  $M$ .

The admittance matrix for the half-space region is identical to that for the problem of electromagnetic transmission through a rectangular aperture in a conducting plane [4]. The admittance

matrix for the waveguide region is of the form of a summation over all modes. The excitation may be taken to be any mode, but normally it is taken to be the dominant mode.

The main program computes the complex coefficients  $V_i$  which determine the magnetic current  $M$  according to

$$M = \sum_i V_i M_i \quad (1)$$

where  $M_i$  are the above described expansion functions, the amplitudes  $\Gamma_i$  of the reflected waveguide modes, the equivalent aperture admittance according to

$$Y = \frac{1 - \Gamma_0}{1 + \Gamma_0} Y_0 \quad (2)$$

the complex power flowing through the aperture, and four radiation gain patterns. The four gain patterns are those of two orthogonal components of  $E$  in the two principal planes.

The equivalent magnetic current in the aperture and the radiation gain patterns for a slot  $\lambda$  by  $\lambda/10$  are shown in Fig. 2. The slot is centered in the cross section of a rectangular waveguide of dimensions  $\lambda$  by  $\lambda/2$  excited in the dominant  $TE_{10}$  mode. The gain patterns are the two principal plane patterns,  $G_{xx}$  denoting an  $x$ -polarized field in the  $x = 0$  plane, and  $G_{\theta y}$  denoting a  $\theta$ -polarized field in the  $y = 0$  plane. Note that the  $M_x$ , or, equivalently, the  $E_y$  in the slot, is nearly the same as the dominant mode field, which is a commonly made assumption.

Fig. 3 shows a comparison of some computed values of the equivalent aperture admittance seen by a waveguide to some measured values and to some computed values from another solution [5]. The waveguide cross section is square and opens into half space through a square aperture of the same size. Other examples of computation of magnetic current, radiation gain patterns, and waveguide admittances for various cases are given in the original report [1].

## REFERENCES

- [1] J. R. Mautz and R. F. Harrington, "Transmission from a rectangular waveguide into half space through a rectangular aperture," Report TR-76-5, Department of Electrical and Computer Engineering, Syracuse University, May 1976
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